

Geometrically exact beam theory

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Introduction

Cosserat beam [1]

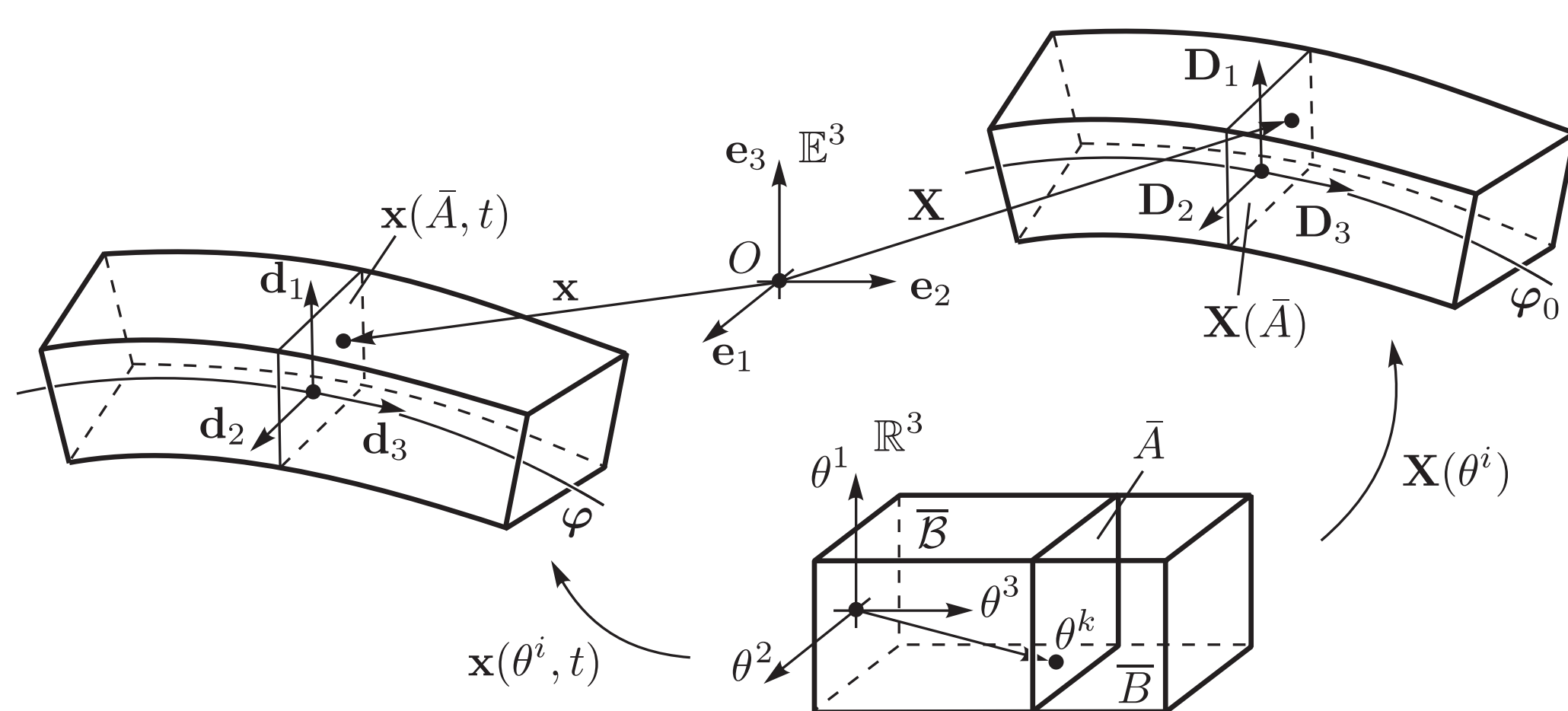
- Nonlinear beam finite elements
- Interpolation of the director field
- Does not rely on rotational degrees of freedom
- Frame indifference and conservation of angular momentum

Director-based theory in skew coordinates [2]

- Formulation in convected coordinates
- Accounts for the lack of orthonormality of the discrete director frame
- Improves dramatically the numerical performance

Geometrically exact beam – Kinematical assumptions

- Reference and current configuration



- Restricted position field

$$\mathbf{x}(\theta^\alpha, s, t) = \varphi(s, t) + \theta^\alpha \mathbf{d}_\alpha(s, t)$$

where $\mathbf{R}(s, t) \in SO(3)$ such that

$$\mathbf{d}_k(s) = \mathbf{R}(s, t) \mathbf{e}_k, \text{ with } \mathbf{R} = \mathbf{d}_k \otimes \mathbf{e}^k$$

- Skew-symmetric effective curvature $\tilde{\mathbf{k}}$

$$(\mathbf{d}_k)_{,s} = \tilde{\mathbf{k}} \mathbf{d}_k = \mathbf{k} \times \mathbf{d}_k, \text{ with } \tilde{\mathbf{k}} = \mathbf{R}_{,s} \mathbf{R}^{-1} = \mathbf{d}_{j,s} \otimes \mathbf{d}^j$$

Virtual work and director formulation

- Virtual work

$$\delta W = \int_{s_1}^{s_2} \left\{ \delta \varphi \cdot (A_p \ddot{\varphi} + \ddot{\mathbf{q}} - \bar{\mathbf{n}} - \mathbf{n}_{,s}) + \delta \phi \cdot (\mathbf{q} \times \ddot{\varphi} + \dot{\mathbf{h}} - \bar{\mathbf{m}} - \mathbf{m}_{,s} - \varphi_{,s} \times \mathbf{n}) \right\} ds + (\delta \varphi \cdot (\mathbf{n} - \bar{\mathbf{n}}) + \delta \phi \cdot (\bar{\mathbf{m}} - \mathbf{m})) \Big|_0^L = 0 \quad \forall \delta \varphi, \delta \phi, t$$

- Constitutive law

$$W(\gamma_i, k^l) = \frac{1}{2} d^{1/2} \gamma_i (\hat{\mathbf{D}}_1)^{ij} \gamma_j + \frac{1}{2} d^{1/2} k^l (\hat{\mathbf{D}}_2)_{ij} k^j$$

with

$$[\hat{\mathbf{D}}_1] = \text{Diag}[GA_1, GA_2, EA] \quad \text{and} \quad [\hat{\mathbf{D}}_2] = \text{Diag}[Eh, Eh, Gj]$$

- Variational formulation of the geometrically exact beam

$$\delta W = \int_{s_1}^{s_2} \left\{ \delta \varphi \cdot [A_p \ddot{\varphi} + q_p^\beta \ddot{\mathbf{d}}_\beta - \bar{\mathbf{n}}] + \delta \mathbf{d}_\alpha \cdot [M_p^{\alpha\beta} \ddot{\mathbf{d}}_\beta + q_p^\alpha \ddot{\varphi}] - \delta \mathbf{d}_i \cdot \frac{1}{2} (\mathbf{d}^i \times \bar{\mathbf{m}}) + \lambda^j \delta \mathbf{d}_i \cdot \mathbf{d}_j + \mathbf{n} \cdot \delta \varphi_{,s} + \frac{1}{2} [(\mathbf{n} \cdot \mathbf{d}^i)(\varphi_{,s} \cdot \delta \mathbf{d}_i) - (\delta \mathbf{d}_i \cdot \mathbf{n})(\mathbf{d}^i \cdot \varphi_{,s})] + \frac{1}{2} (\mathbf{m} \cdot \mathbf{d}_i) \varepsilon_{ijk} [\delta (d^{-1/2}) (\mathbf{d}_k \cdot \mathbf{d}_{j,s}) + d^{-1/2} \delta \mathbf{d}_k \cdot \mathbf{d}_{j,s} + d^{-1/2} (\mathbf{d}_k \cdot \delta \mathbf{d}_{j,s})] \right\} ds - \left(\delta \varphi \cdot \bar{\mathbf{n}} + \delta \mathbf{d}_i \cdot \frac{1}{2} (\mathbf{d}^i \times \bar{\mathbf{m}}) \right) \Big|_{s_1}^{s_2} = 0 \quad \forall \delta \varphi, \delta \mathbf{d}_i, t$$

Finite element formulation

- Galerkin type approach

$$\varphi^h(s, t) = \sum_A N^A(s) \varphi_A(t), \quad \mathbf{d}_i^h(s, t) = \sum_A N^A(s) \mathbf{d}_{iA}(t)$$

$$\delta \varphi^h(s) = \sum_A N^A(s) \delta \varphi_A, \quad \delta \mathbf{d}_i^h(s) = \sum_A N^A(s) \delta \mathbf{d}_{iA}$$

- Virtual work of the contact force

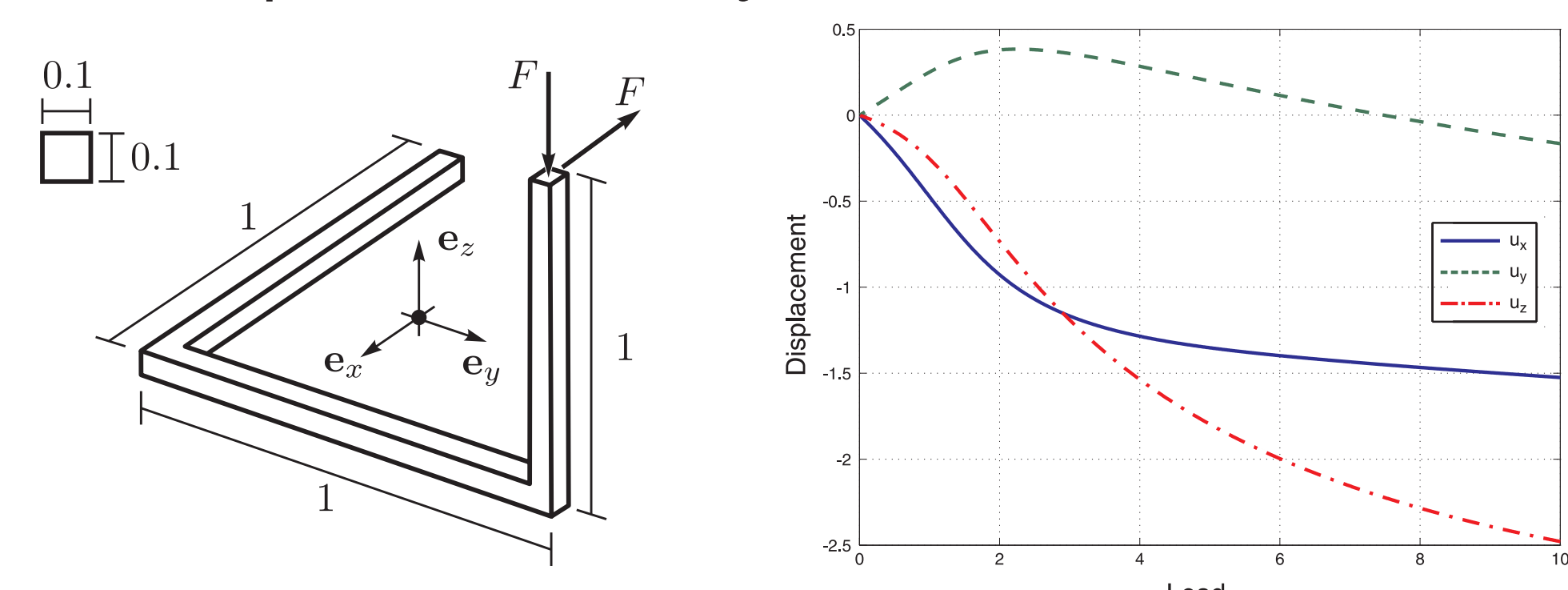
$$\delta W_1^{\text{int,h}} = \delta \varphi_A \cdot \int_{s_1}^{s_2} N_{,s}^A \mathbf{n}^h ds + \sum_{i=1}^3 \delta \mathbf{d}_{iA} \cdot \int_{s_1}^{s_2} \frac{1}{2} [(\mathbf{n}^h \cdot \mathbf{d}^{i,h}) N^A N_{,s}^B \varphi_B - (\varphi_{,s}^h \cdot \mathbf{d}^{i,h}) N^A \mathbf{n}^h] ds$$

- Virtual work of the contact torques

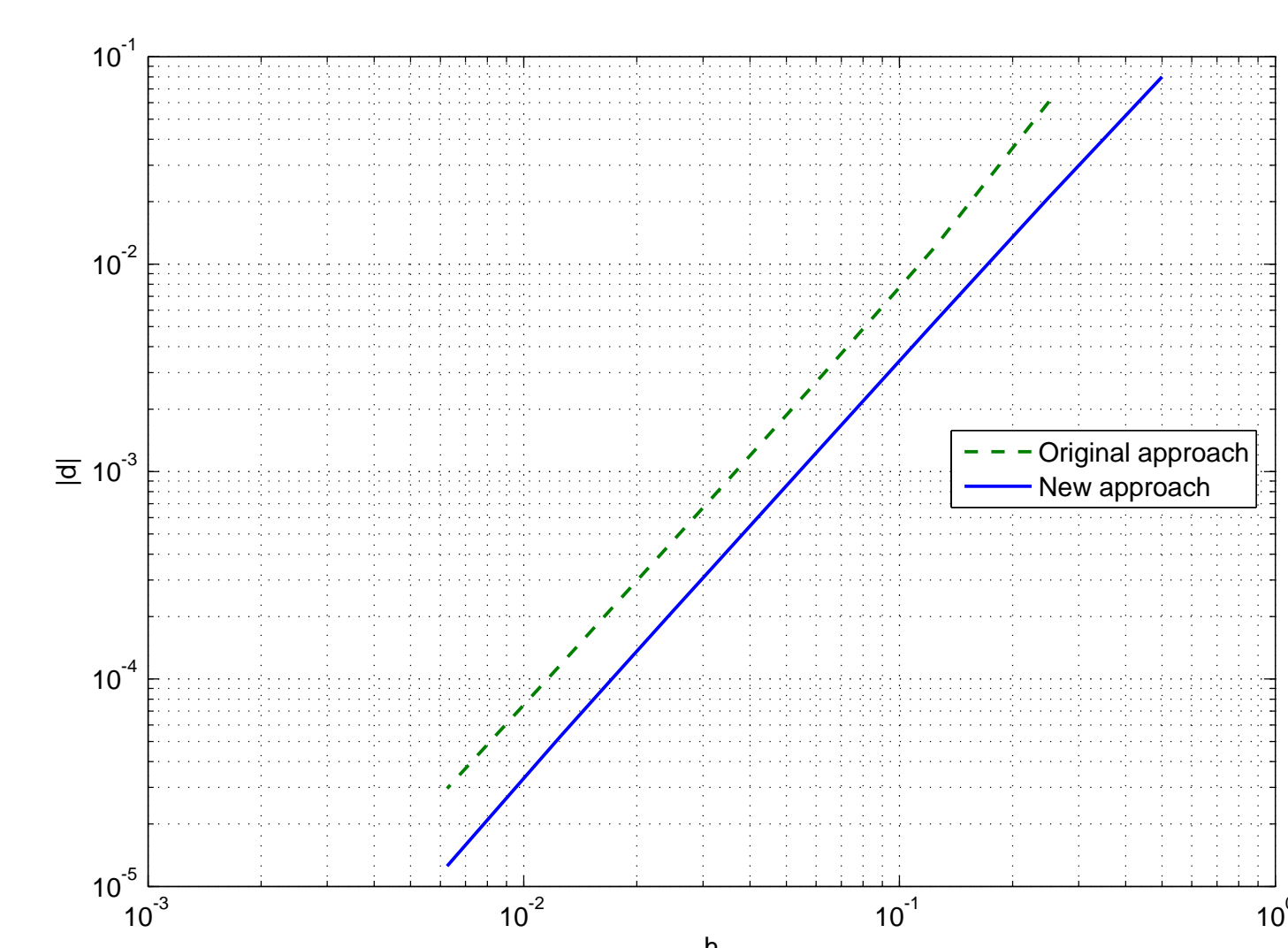
$$\delta W_2^{\text{int,h}} = \frac{1}{2} \delta \mathbf{d}_{iA} \cdot \int_{s_1}^{s_2} \det(\mathbf{R}^{-1,h}) (\mathbf{m}^h \cdot \mathbf{d}_k^h) \varepsilon_{kij} (N^A N_{,s}^B \mathbf{d}_{jB} - N_{,s}^A N^B \mathbf{d}_{jB}) ds - \frac{1}{2} \delta \mathbf{d}_{1A} \cdot \int_{s_1}^{s_2} (\det(\mathbf{R}^{-1,h}))^2 N^A (\mathbf{d}_2^h \times \mathbf{d}_3^h) (\mathbf{m}^h \cdot \mathbf{d}_1^h) \varepsilon_{lmn} (\mathbf{d}_n^h \cdot \mathbf{d}_{m,s}^h) ds - \frac{1}{2} \delta \mathbf{d}_{2A} \cdot \int_{s_1}^{s_2} (\det(\mathbf{R}^{-1,h}))^2 N^A (\mathbf{d}_3^h \times \mathbf{d}_1^h) (\mathbf{m}^h \cdot \mathbf{d}_2^h) \varepsilon_{lmn} (\mathbf{d}_n^h \cdot \mathbf{d}_{m,s}^h) ds - \frac{1}{2} \delta \mathbf{d}_{3A} \cdot \int_{s_1}^{s_2} (\det(\mathbf{R}^{-1,h}))^2 N^A (\mathbf{d}_1^h \times \mathbf{d}_2^h) (\mathbf{m}^h \cdot \mathbf{d}_3^h) \varepsilon_{lmn} (\mathbf{d}_n^h \cdot \mathbf{d}_{m,s}^h) ds$$

Numerical example

Beam with slope discontinuity



Geometry of the structure and tip displacement versus load F .



Errors in tip displacements versus h-refinement.

References

- P. Betsch and P. Steinmann
Frame-indifferent beam finite elements based upon the geometrically exact beam theory.
Int. J. Numer. Meth. Engng, 54:1775–1788, 2002.
- S.R. Eugster, C. Hesch, P. Betsch and Ch. Glocker
Director-based beam finite elements relying on the geometrically exact beam theory formulated in skew coordinates
Int. J. Numer. Meth. Engng, 97:111–129, 2014